

ABSORPTION OF HIGH-ENERGY PHOTONS IN THE UNIVERSE

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The probability per unit length of path that a  $10^{12}$ -ev  $\gamma$  quantum is converted into an electron pair as a result of a collision with a thermal photon is calculated. If the energy density of thermal photons in intergalactic space is taken as  $0.1 \text{ ev cm}^{-3}$ , the probability turns out to be  $7 \times 10^{-27}$ . Thus if the distance traversed is greater than  $10^{26}$  cm, the attenuation of the  $\gamma$ -quantum flux may be appreciable.

THERE has recently been increasing interest in the possibility of observing point sources of high energy photons.<sup>[1]</sup> In this article, we shall consider the role of the reaction  $\gamma + \gamma \rightarrow e^+ + e^-$  in the propagation of  $10^{12} - 10^{13}$ -ev photons from sufficiently distant objects outside our galaxy.

The cross section for the conversion of two  $\gamma$  quanta into an electron pair is given by the expression\* (see [2])

$$\sigma(s) = \frac{1}{2} \pi r_0^2 (1 - v^2) \left\{ (3 - v^4) \ln \frac{1+v}{1-v} + 2v(v^2 - 2) \right\},$$

$$v = \sqrt{1 - 1/s},$$

$$r_0 = 2.8 \cdot 10^{-13} \text{ cm}, s = (E\epsilon/2m^2) (1 - \cos \theta),$$

where  $s$  is the square of the c.m.s.  $\gamma$ -quantum energy,  $m$  is the mass of the electron,  $c = 1$ ,  $E$  and  $\epsilon$  are the energies of the colliding  $\gamma$  quanta in the laboratory system,  $\theta$  is the angle between their momenta;  $\sigma(s) \approx 10^{-25} \text{ cm}^2$  in the region of  $s$  of interest to us. At present, it is assumed that the density of photons with mean energy  $\sim 1 \text{ ev}$  in intergalactic space is  $1/3$  to  $1/10$  the density in our galaxy. The density of light energy in the galaxy is  $W_{\text{gal}} = 0.3 - 1 \text{ ev}$ .<sup>[3]</sup> It is thus readily seen that if the path traversed by high energy photons is  $R \gtrsim 10^{26} \text{ cm}$ , then the photon flux can be appreciably attenuated. Similar estimates indicate that the contribution to the attenuation of the photon beam as a result of interactions with nuclei or magnetic fields is much smaller.

We proceed to quantitative estimates. The probability per unit length of path that a quantum of en-

\*It is readily seen that  $\sigma(s)$  is obtained by multiplication of the inverse reaction by  $2v^2$ ; the factor 2 results from the fact that the particles in the final state are not identical and  $v^2$  results from the difference in the flux and statistical weight of these channels of the reaction.

ergy  $E$  is converted into an electron pair in a collision with a thermal photon is

$$P = 2 \int_0^\infty d\epsilon n(\epsilon) \int_0^1 z \sigma(s) dz, z = \frac{1}{2} (1 - \cos \theta),$$

$n(\epsilon)$  is the density of thermal photons in the energy interval  $d\epsilon$ . Replacing the integration over  $z$  by integration over  $s = E\epsilon z/m^2$ , we find that

$$P = 2 \left( \frac{m^2}{E} \right)^2 \int_0^\infty n(\epsilon) \epsilon^{-2} \varphi(s_0) d\epsilon, \varphi(s_0) = \int_1^{s_0} \sigma(s) ds, s_0 = \frac{E\epsilon}{m^2}.$$

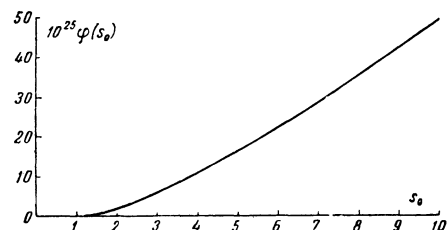
The values of  $\varphi(s_0)$  in the interval  $1 \leq s_0 \leq 10$  are shown in the figure. For larger  $s_0$

$$\varphi(s_0) = 25 \cdot 10^{-26} \{s_0 (\ln 4s_0 - 2) + 3\}.$$

For the numerical estimate, we set

$$n(\epsilon) = A \epsilon^2 / (e^{2\epsilon} - 1).$$

This is a spectrum of the solar type, where  $kT = 0.5$  and the photon energy is measured in electron-volts. To consider a specific case, we shall assume that the energy density of thermal photons in the universe is  $0.1 \text{ ev cm}^{-3}$ . Then the normalization factor is  $A = 0.22$ . Shown in the table are the numerical values of  $P$  for different  $\gamma$ -quantum energies and, as an example, the values of  $PR$  for an interesting star, Cygnus A (at a distance<sup>[4]</sup>  $R_C = 6.6 \times 10^{26} \text{ cm}$ ). It is seen from the table that



$10^{-12} E, \text{ ev}$	0.1	0.5	1	5	10	50
$10^{27} P, \text{ cm}^{-1}$	0.05	5	7	4	2	0.7
$PR_c$	0.03	3	4.6	2.6	1.3	0.5

the maximum attenuation of the beam is  $e^{-PR}$  for  $E = 10^{12}$  ev.

In principle, the effect can be used for an experimental estimate of the mean density of thermal photons in intergalactic space. The numerical value of this density is of interest for a number of astrophysical problems (see, e.g., <sup>[5]</sup>, where the photo-disintegration of high energy heavy nuclei in intergalactic space is discussed).

In conclusion, the author expresses his gratitude to V. L. Ginzburg for interesting discussions.

<sup>1</sup>G. Cocconi, Proceedings of the Moscow Cosmic Ray Conference, 1960, vol. 2, p. 309; Sekido, Yoshida, Komiya, Heno, and Murayama, *ibid.*, vol. 3, pp. 137, 140; M. P. Savedoff, *Nuovo cimento* **13**, 12 (1959); P. Morrison, *Nuovo cimento* **7**, 858 (1958).

<sup>2</sup>A. I. Akiezer and V. B. Berestetskii, *Kvantovaya élektrodinamika (Quantum Electrodynamics)*, 2nd ed., Fizmatgiz, 1959, p. 359.

<sup>3</sup>E. Feenberg and H. Primakoff, *Phys. Rev.* **73**, 449 (1948); C. W. Allen, *Astrophysical Quantities*, University of London, Athlone Press, 1955, pp. 228, 245.

<sup>4</sup>I. S. Shklovskii, *Astronomicheskii zhurnal*, **37**, 945 (1960), *Soviet Astronomy* **4**, 885 (1961).

<sup>5</sup>N. M. Gerasimova and I. L. Rozental', *JETP* **41**, 488 (1961), *Soviet Phys. JETP* **14**, 350 (1962).