

INVESTIGATION OF THE ANGULAR DISTRIBUTIONS OF  $\alpha$  PARTICLES FROM THE  $\text{Li}^7(p, \alpha)\text{He}^4$  REACTION

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The angular distributions of  $\alpha$  particles emitted in the  $\text{Li}^7(p, \alpha)\text{He}^4$  reaction at angles between 20 and 160° were measured for incident proton energies of 5.78, 6.15, and 6.55 MeV. The excitation curve for proton energies between 3.3 and 6.6 MeV was found to have a distinct resonance structure. The results are analyzed.

THE angular distributions of  $\alpha$  particles produced in the  $\text{Li}^7(p, \alpha)\text{He}^4$  reaction have been studied<sup>[1]</sup> only for proton energies  $E_p < 3.75$  MeV. These measurements indicate the existence of an excitation resonance on the excitation curve at  $E_p \sim 3.0$  MeV. An analysis of the angular distributions in this energy region has been made by Inglis.<sup>[2]</sup>

In the present experiment, the  $\text{Li}^7(p, \alpha)\text{He}^4$  reaction has been investigated for proton energies in the interval 3.3 – 6.55 MeV. The  $\alpha$  particles produced in the reaction were recorded by a "telescope" consisting of three proportional counters.<sup>[3]</sup> Targets of  $\text{Li}_2\text{CO}_3$  deposited on gold tinsel by evaporation in vacuo were used in the measurements. The target thicknesses were 0.16, 0.52, and 0.92 mg/cm<sup>2</sup>, which corresponded to losses of 13, 41, and 70 keV for 6-MeV protons.

Figure 1 shows the range distribution of the  $\alpha$  particles and  $\text{He}^3$  nuclei emitted at 30° from the target in the case of 6.5-MeV incident protons. The expected positions of the  $\text{He}^4$  and  $\text{He}^3$  groups formed in the  $\text{Li}^7(p, \alpha)\text{He}^4$  and  $\text{Li}^6(p, \alpha)\text{He}^3$  reactions are indicated by arrows.

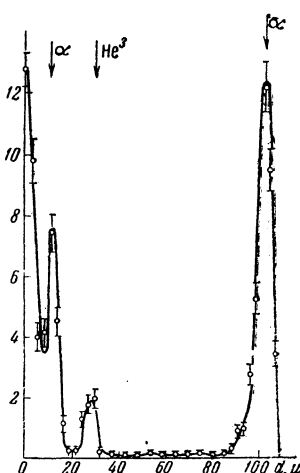


FIG. 1. Range distribution of  $\alpha$  particles and  $\text{He}^3$  nuclei emitted at an angle of 30° in the  $\text{Li}(p, \alpha)$  reaction ( $E_p = 6.5$  MeV). The abscissa axis gives the thickness of an aluminum absorber in front of the first counter.

The angular distributions of the  $\alpha$  particles produced in the  $\text{Li}^7(p, \alpha)\text{He}^4$  reaction ( $Q = 17.35$  MeV) were measured for three incident proton energies:  $E_p = 5.78, 6.15,$  and  $6.55$  MeV. The results were recalculated for the c.m.s. and are shown in Figs. 2a, b, c. Since two identical particles are produced in this reaction, the angular distributions in this reaction can be represented in the form of a series of even powers of  $\cos \vartheta$  (where  $\vartheta$  is the c.m.s. angle) or in the form of an expansion in even Legendre polynomials  $P_n(x)$ . From the experimental distributions we calculated, by the method of least squares, the coefficients in expressions of the form

$$d\sigma/d\Omega = (\sigma/4\pi) [1 + a_2P_2(x)],$$

$$d\sigma/d\Omega = (\sigma/4\pi) [1 + A_2P_2(x) + A_4P_4(x)]$$

where  $\sigma$  is the total cross section.

We use the following values for the coefficients

$E_p, \text{MeV}$	$a_2$	$A_2$	$A_4$
6.55	-0.169	-0.167	-0.102
6.15	-0.357	-0.356	-0.010
5.78	-0.717	-0.693	+0.085

The results of the calculation of the angular distributions from the foregoing formulas are shown in Figs. 2a, b, c [the dashed curves correspond to the formula with two terms and the solid curves to the formula with  $P_4(x)$ ]. It is seen from the figures that even the simple formula of the form  $1 + a_2P_2(x)$  satisfactorily describes the experimental results.

The excitation curves were measured in the proton energy regions 5.25 – 6.55 MeV at 30° (35° in the c.m.s.) and 3.3 – 6.55 MeV at 80° (90° in the c.m.s.). The results of the measurements are shown in Fig. 3. Also shown in the same figure are the data for 90° obtained by other authors (see<sup>[1]</sup>) ( $E_p = 1.0 - 3.75$  MeV). Since in our and in other experiments<sup>[1]</sup> only the relative cross sec-

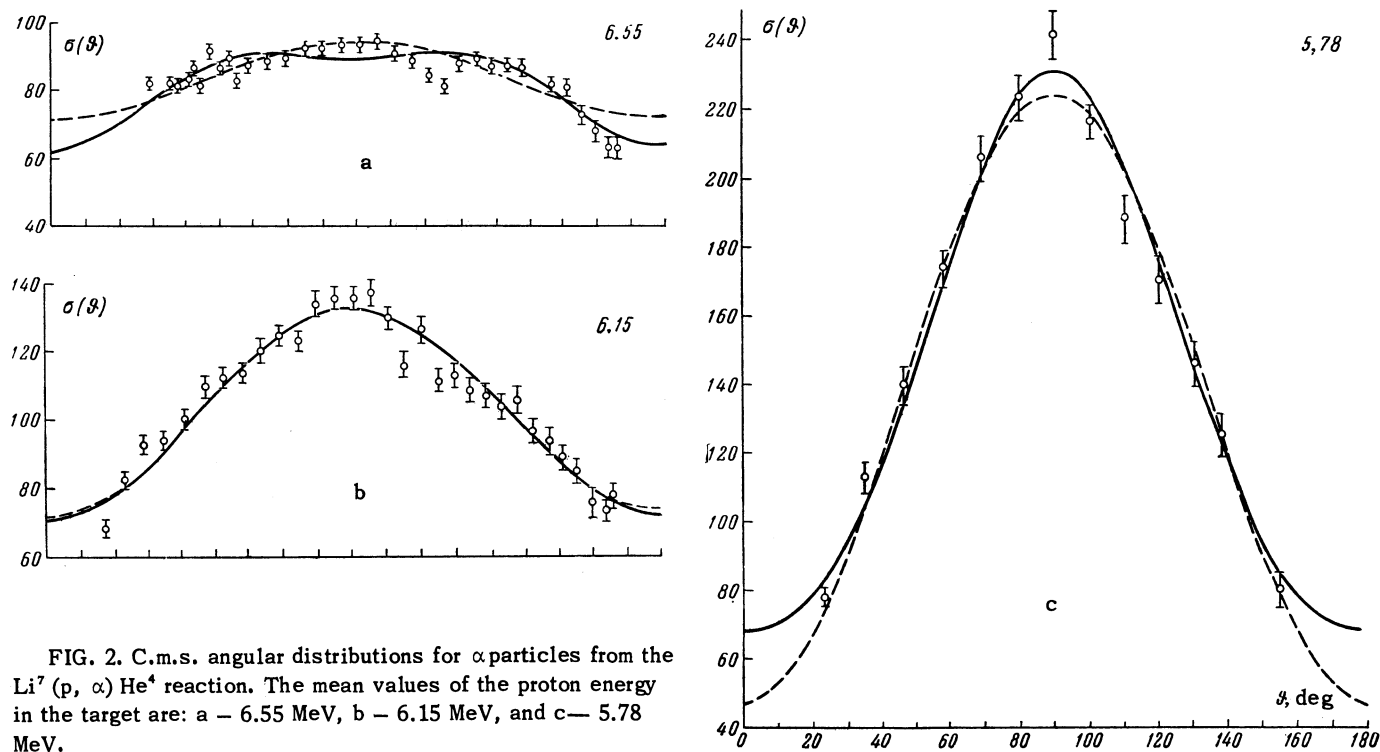


FIG. 2. C.m.s. angular distributions for  $\alpha$  particles from the  $\text{Li}^7(p, \alpha)\text{He}^4$  reaction. The mean values of the proton energy in the target are: a — 6.55 MeV, b — 6.15 MeV, and c — 5.78 MeV.

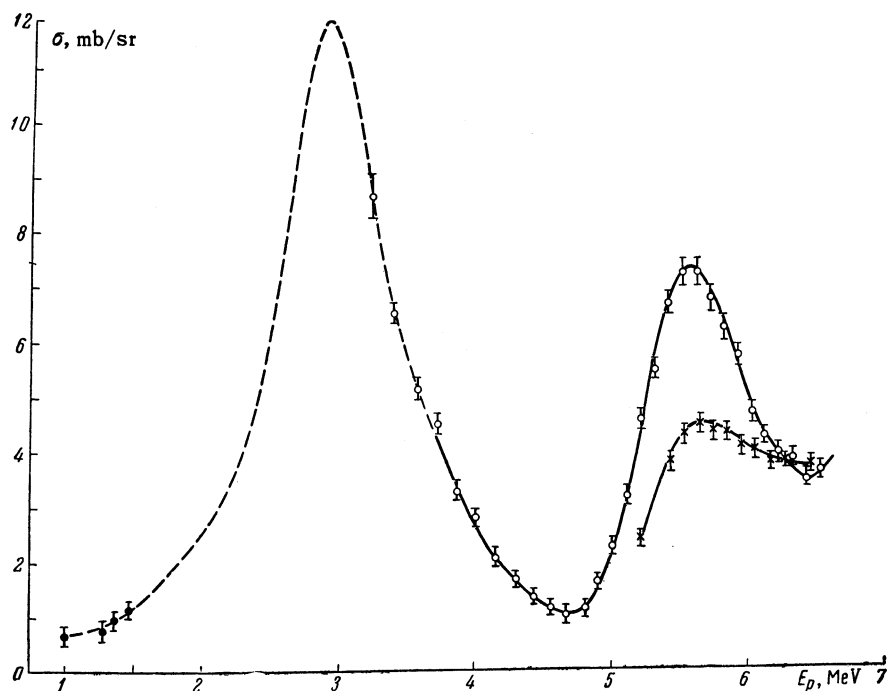


FIG. 3. Excitation energy as a function of the proton energy:  $\times$  — for  $\alpha$  particles emitted at  $30^\circ$  l.s.;  $\circ$  — at  $80^\circ$  l.s.;  $\bullet$  — data of Freeman et al.<sup>[4]</sup> The results obtained by Heydenburg et al.<sup>[1]</sup> are shown dashed.

tions were measured, the excitation curves for  $80^\circ$  and  $90^\circ$  were fitted together in the overlapping region (3.3 — 3.75 MeV). It is seen from Fig. 3 that the curves fit together well; here, the effect due to the difference in the angles of measurement ( $80^\circ$  in our experiment and  $90^\circ$  in the others) is at most 3%. In order to determine the absolute differential cross sections given on the ordinate axis, we used

the results of Freeman, Hanna, and Montague<sup>[4]</sup> at  $90^\circ$  for several values of  $E_p$ .

From the data for  $30^\circ$  and  $80^\circ$  we calculated the values of the coefficient  $a_2$  for proton energies between 5.25 and 6.5 MeV. The obtained values (black circles in Fig. 4) indicate that  $a_2 < 0$  and that  $|a_2|$  increases as the resonance value of  $E_p$  is approached. Shown in the same figure (empty

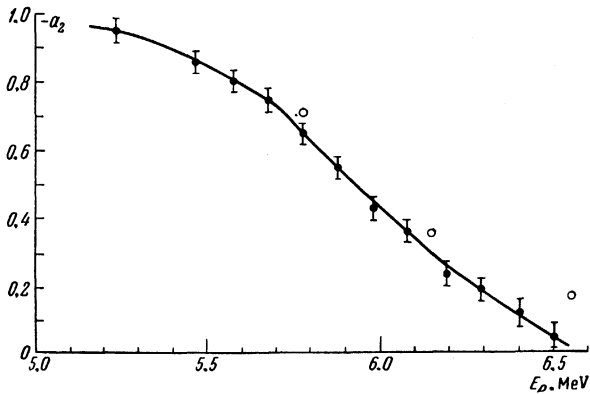


FIG. 4. The coefficient  $a_2$  as a function of the proton energy.

circles) are the values of the coefficient  $a_2$  calculated by the method of least squares from the angular distributions (see the data given above).

The quite distinct resonance structure of the excitation curve for the  $\text{Li}^7(p, \alpha)\text{He}^4$  reaction for incident proton energies up to 6.6 MeV (Fig. 3) indicates that the basic mechanism for the reaction is a mechanism associated with the formation of a  $\text{Be}^8$  compound nucleus. This is also confirmed by the application of the expressions  $1 + a_2P_2(x)$  and  $1 + A_2P_2(x) + A_4P_4(x)$  to describe the angular distributions.

The excitation curve (Fig. 3) contains two resonances: one in the region 3.0 MeV and the other in the region 5.6 MeV. The resonance at 3 MeV was analyzed by Inglis,<sup>[2]</sup> who concluded that the angular distributions of the  $\alpha$  particles in the region of this resonance can be explained if it is assumed that the  $\text{Be}^8$  compound nucleus has a  $2^+$  state with excitation energy 19.9 MeV and a width  $\Gamma = 1$  MeV and a  $0^+$  state with a width of several MeV lying above the resonance region.

The existence of the resonance at 5.6 MeV is most simply explained by the occurrence of a state with excitation energy 22.3 MeV and width  $\Gamma = 1$  MeV in the  $\text{Be}^8$  nucleus. In a review article, Ajzenberg-Selove and Lauritsen<sup>[5]</sup> listed a level with excitation energy 22.6 MeV observed by Whaling and Bonner.<sup>[6]</sup> These latter authors, who investigated the excitation curves for the  $\text{Li}^6(d, \alpha)\text{He}^4$  and  $\text{Li}^6(d, p)\text{Li}^7$  reactions found that the function  $\alpha(E_d) = \sigma(90^\circ)/\lambda^2 P_0(E_d)$  (where  $P_0$  is the penetrability for  $s$  deuterons) has a resonance at  $E_d = 0.4$  MeV. However, Hirst, Johnstone, and Poole<sup>[7]</sup> did not observe this level in their study of the  $\text{Li}^6(d, \alpha)\text{He}^4$  and  $\text{Li}^6(d, n)\text{Be}^7$  reactions. Comparison of these two experiments indicates that the results of Hirst et al.<sup>[7]</sup> are apparently more reliable. First, in the latter work, the authors used tables of Coulomb functions in the calculation of the penetrability  $P_0$ , while Whaling and

Bonner<sup>[6]</sup> calculated  $P_0$  with the aid of approximation formulas. Second, Hirst et al.<sup>[7]</sup> used a thin target (0.005 mg/cm<sup>2</sup>), while Whaling and Bonner<sup>[6]</sup> employed a comparatively thick target (0.15 mg/cm<sup>2</sup>), which could have affected the results of the cross section measurements for low-energy deuterons. Hence the existence of a 22.6-MeV excited state in  $\text{Be}^8$  cannot be regarded as established.

Thus, it is most natural to assume that the resonance at  $E_p = 5.6$  MeV is connected with a single level of the compound nucleus having an excitation energy of 22.3 MeV and width  $\Gamma = 1$  MeV. In determining the characteristics of this level we should keep in mind the fact that only states of  $\text{Be}^8$  having even spins ( $0^+$ ,  $2^+$ ,  $4^+$ , etc.) can decay into two  $\alpha$  particles. From the law of conservation of angular momentum it follows that the  $\text{Li}^7(p, \alpha)\text{He}^4$  reaction involves only protons with orbital angular momenta  $l = J \pm 1$ . Since, for  $E_p = 6$  MeV, the ratio of the penetrabilities  $P_1 : P_3 : P_5$  is 1 : 0.11 : 0.0006,<sup>[8]</sup> only protons with  $l$  equal to 1 and 3 need to be taken into account.

It thus follows that the possible states are  $0^+$ ,  $2^+$ , and  $4^+$ . The value  $0^+$  drops out since a  $0^+$  state would lead to isotropic angular distributions. The results of the calculation by the method of Inglis<sup>[2]</sup> show that for the  $4^+$  state, with only the  $f$  protons taken into account, the coefficient  $a_2$  should be positive. This result also holds if, in addition to the  $4^+$  state, the  $0^+$  state is present (for example, a broad state whose existence was suggested by Inglis to explain the angular distributions in the region  $E_p = 3$  MeV). Since  $a_2 < 0$  in the region of the 5.6 MeV resonance (see the values of the coefficients and Fig. 4), the value  $4^+$  also drops out. Hence, if the 5.6-MeV resonance is associated only with one state of the compound nucleus with an excitation energy of 22.3 MeV, the spin and parity of this state are  $2^+$ .

It should be mentioned that Gibbons and Macklin,<sup>[9]</sup> who studied the  $\text{Li}^7(p, n)\text{Be}^7$  reaction, observed a resonance in the yield curve of this reaction at  $E_p = 5.0$  MeV. Since we did not find this resonance in the excitation curve measured by us, it is evidently associated with a  $\text{Be}^8$  state which does not decay into two  $\alpha$  particles.

<sup>1</sup> Heydenburg, Hudson, Inglis, and Whitehead, Phys. Rev. **73**, 241 (1948); *ibid.*, **74**, 405 (1948).

<sup>2</sup> D. R. Inglis, Phys. Rev. **74**, 21 (1948).

<sup>3</sup> Teplov, Shevchenko, and Ruuge, JETP **39**, 923 (1960), Soviet Phys. JETP **12**, 640 (1961).

<sup>4</sup> Freeman, Hanna, and Montague, Nuclear Phys. **5**, 148 (1958).

<sup>5</sup>F. Ajzenberg-Selove and T. Lauritsen, Nuclear Phys. **11**, 1 (1959).

<sup>6</sup>W. Whaling and T. W. Bonner, Phys. Rev. **79**, 258 (1950).

<sup>7</sup>Hirst, Johnstone, and Poole, Phil. Mag. **45**, 762 (1954).

<sup>8</sup>Luk'yanov, Teplov, and Akimova, Tablitsy

volnovykh kulonovskikh funktsii (Tables of Coulomb Wave Functions), AN SSSR, 1961.

<sup>9</sup>J. H. Gibbons and R. L. Macklin, Phys. Rev. **114**, 571 (1959).

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